

Optics & Spectroscopy

lecture 13

- ~ There are a number of "model" problems in QM. That can be solved exactly
 - ~ harmonic oscillator
 - ~ rigid rotor
 - ~ H-atom
 - ~ particle-in-a-box
 - ~ most real systems do not perfectly fit these model problems
 - ~ handle this through Perturbation Theory

$$\text{real Hamiltonian} \quad \hat{H} = \hat{H}_0 + \hat{H}'$$

\hat{H}_0
 model
 Hamiltonian

extra lit

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- goal: obtain estimates for the energy levels & wavefunctions of \hat{H}
in terms of our model problem
- ~ do this by expanding energies & wavefunctions in terms of a perturbation parameter " λ "

$$\hat{H} = \hat{H}_0 + \lambda \hat{H}_1$$

$$|\psi_n\rangle = |\psi_n^{(0)}\rangle + \lambda |\psi_n^{(1)}\rangle + \lambda^2 |\psi_n^{(2)}\rangle + \dots$$

$$E_n = E_n^{(0)} + \lambda E_n^{(1)} + \lambda^2 E_n^{(2)} + \dots$$

$|\psi^{(0)}\rangle$ & $E_n^{(0)}$ are the solutions of the model problem

"zeroth-order" solutions

$$\text{i.e. } H_0 |\psi^{(0)}\rangle = E_n^{(0)} |\psi^{(0)}\rangle$$

$|\psi^{(0)}\rangle$ & $E_n^{(0)}$ are known

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$E_n^{(0)}, E_n^{(2)}$

$| \psi_n^{(0)} \rangle, | \psi_n^{(2)} \rangle$

} 1st and 2nd order
corrections to the
wavefunction & energy

write the S.E. $\hat{H} |\psi_n\rangle = E_n |\psi_n\rangle$ as

$$(\hat{H} - E_n) |\psi_n\rangle = 0$$

$$\Rightarrow \left\{ \hat{H}_0 + \lambda \hat{H}' - (E_n^{(0)} + \lambda E_n^{(1)} + \lambda^2 E_n^{(2)} + \dots) \right\}$$

$$\times \left\{ | \psi_n^{(0)} \rangle + \lambda | \psi_n^{(1)} \rangle + \lambda^2 | \psi_n^{(2)} \rangle + \dots \right\}$$

$$= 0$$

expand & collect terms in λ

$$\left\{ \hat{H}_0 - E_n^{(0)} \right\} | \psi_n^{(0)} \rangle$$

$$+ \lambda' \left\{ (\hat{H}_0 - E_n^{(0)}) | \psi_n^{(1)} \rangle + (\hat{H}' - E_n^{(1)}) | \psi_n^{(0)} \rangle \right\}$$

$$+ \lambda^2 \left\{ (\hat{H}_0 - E_n^{(0)}) | \psi_n^{(2)} \rangle + (\hat{H}' - E_n^{(1)}) | \psi_n^{(1)} \rangle + E_n^{(2)} | \psi_n^{(0)} \rangle \right\}$$

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$$+ \lambda^3 \{ \} + \lambda^4 \{ \} + \dots = 0$$

\sim now $\lambda \neq 0$ ($\lambda = 1$ \sim bookkeeping device)

\Rightarrow each term that multiplies λ must be separately equal to zero

$$\Rightarrow (\hat{H}^{(0)} - E_n^{(0)}) |\psi_n^{(0)}\rangle = 0 \quad \begin{matrix} \leftarrow & \text{model} \\ & \text{problem} \end{matrix}$$

$$\lambda' : (\hat{H}_0 - E_n^{(0)}) |\psi_n^{(1)}\rangle + (\hat{H}' - E_n^{(1)}) |\psi_n^{(0)}\rangle = 0$$

$$\lambda'' : (\hat{H}_0 - E_n^{(0)}) |\psi_n^{(0)}\rangle + (\hat{H}' - E_n^{(1)}) |\psi_n^{(1)}\rangle - E_n^{(2)} |\psi_n^{(0)}\rangle = 0$$

1st-order P.T.

$$(\hat{H}_0 - E_n^{(0)}) |\psi_n^{(0)}\rangle + (\hat{H}' - E_n^{(1)}) |\psi_n^{(0)}\rangle = 0$$

$\overbrace{\text{know } |\psi_n^{(0)}\rangle}$ \sim complete basis set

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\Rightarrow expand $|\psi_n^{(1)}\rangle$ in terms of $\{|\psi_i^{(0)}\rangle\}$

$$|\psi_n^{(1)}\rangle = \sum_i c_{ni} |\psi_i^{(0)}\rangle$$

$$c_{ni} = \langle \psi_i^{(0)} | \psi_n^{(1)} \rangle$$

$$\Rightarrow (\hat{H}_0 - E_n^{(0)}) \sum_i c_{ni} |\psi_i^{(0)}\rangle + (\hat{H}' - E_n^{(1)}) |\psi_n^{(0)}\rangle = 0$$

~ multiply from the left by $\langle \psi_n^{(0)} |$

$$\langle \psi_n^{(0)} | (\hat{H}_0 - E_n^{(0)}) \sum_i c_{ni} |\psi_i^{(0)}\rangle$$

$$= - \langle \psi_n^{(0)} | (\hat{H}' - E_n^{(1)}) | \psi_n^{(0)} \rangle$$

$$\langle \psi_n^{(0)} | (E_n^{(0)} - E_n^{(0)}) \sum_i c_{ni} |\psi_i^{(0)}\rangle$$

$$= - \langle \psi_n^{(0)} | \hat{H}' | \psi_n^{(0)} \rangle + E_n^{(1)} \underbrace{\langle \psi_n^{(0)} | \psi_n^{(0)} \rangle}_{=1}$$

$$\text{LHS} = 0$$

$$\Rightarrow E_n^{(1)} = \langle \psi_n^{(0)} | \hat{H}' | \psi_n^{(0)} \rangle$$

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~ 1st order correction to the energy
 is just an average of the
 perturbation over the zeroth-order
 wavefunctions.

$$\text{Back to } (\hat{H}_0 - E_n^{(0)}) |\psi_n^{(1)}\rangle = -(\hat{H}' - E_n^{(1)}) |\psi_n^{(0)}\rangle$$

$$\text{~again expand } |\psi_n^{(1)}\rangle = \sum_i C_{ni} |\psi_i^{(0)}\rangle$$

$$(\hat{H}_0 - E_n^{(0)}) \sum_i C_{ni} |\psi_i^{(0)}\rangle = -(\hat{H}' - E_n^{(1)}) |\psi_n^{(0)}\rangle$$

~ now multiply from LHS by $\langle \psi_j^{(0)} | j \neq n$

$$\langle \psi_j^{(0)} | (\hat{H}_0 - E_n^{(0)}) \sum_i C_{ni} |\psi_i^{(0)}\rangle$$

$$= - \langle \psi_j^{(0)} | (\hat{H}' - E_n^{(0)}) |\psi_n^{(0)}\rangle$$

$$\text{LHS} = \sum_i C_{ni} (E_j^{(0)} - E_n^{(0)}) \langle \psi_j^{(0)} | \psi_i^{(0)} \rangle$$

$$= \sum_i C_{ni} (E_j^{(0)} - E_n^{(0)}) \delta_{ji}$$

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~ perform the sum over i , we pick out the term $i=j$

$$\text{LHS} = C_{nj} (E_j^{(0)} - E_n^{(0)})$$

$$\begin{aligned}\text{RHS} &= - \langle \psi_j^{(0)} | \hat{H}' | \psi_n^{(0)} \rangle + E_n^{(0)} \underbrace{\langle \psi_j^{(0)} | \psi_n^{(0)} \rangle}_{\substack{=0 \\ j \neq n}} \\ &= - \langle \psi_j^{(0)} | \hat{H}' | \psi_n^{(0)} \rangle\end{aligned}$$

$$\Rightarrow C_{nj} (E_j^{(0)} - E_n^{(0)}) = - \langle \psi_j^{(0)} | \hat{H}' | \psi_n^{(0)} \rangle$$

$$\Rightarrow C_{nj} = - \frac{\langle \psi_j^{(0)} | \hat{H}' | \psi_n^{(0)} \rangle}{E_j^{(0)} - E_n^{(0)}}$$

$$\Rightarrow |\psi_n\rangle \simeq |\psi_n^{(0)}\rangle + |\psi_n^{(1)}\rangle$$

$$= |\psi_n^{(0)}\rangle + \sum_i C_{ni} |\psi_i^{(0)}\rangle$$

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$$\Rightarrow |\psi_n\rangle \approx |\psi_n^{(0)}\rangle + \sum_{i \neq n} -\frac{\langle \psi_i^{(0)} | \hat{H}' | \psi_n^{(0)} \rangle}{E_i^{(0)} - E_n^{(0)}} |\psi_i^{(0)}\rangle$$

wave in a more compact notation:

$$|\psi\rangle \approx |\psi^0\rangle + \sum_{i \neq n} \underbrace{\frac{\langle i^0 | \hat{H}' | \psi^0 \rangle}{E_n^0 - E_i^0}}_{\text{1st-order correction}} |i^0\rangle$$

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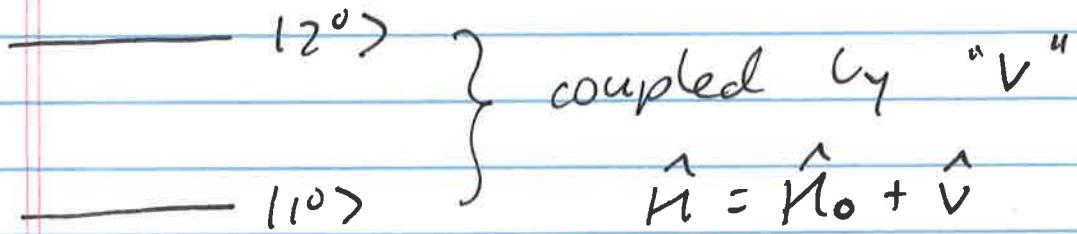
1<sup>st</sup>-order correction  
to the wavefunction

~ perturbation ( $\hat{H}'$ ) mixes other states  
into the wavefunction

~ degree of mixing depends on coupling  
( $\langle i^0 | \hat{H}' | \psi^0 \rangle$ ) and how close the  
2 states are in energy ( $\frac{1}{E_n^{(0)} - E_i^{(0)}}$ )

(a)

## Two-level System:



~ define  $\Delta = E_2^o - E_1^o$   
 $V_{12} = \langle 1^o | \hat{V} | 2^o \rangle$

## 1<sup>st</sup>-order Perturbation Theory

$$|2\rangle \simeq |2^o\rangle + \frac{\langle 1^o | \hat{V} | 2^o \rangle}{E_2^o - E_1^o} |1^o\rangle = |2^o\rangle + \frac{V_{12}}{\Delta} |1^o\rangle$$

$$|1\rangle \simeq |1^o\rangle + \frac{\langle 2^o | \hat{V} | 1^o \rangle}{E_1^o - E_2^o} |2^o\rangle = |1^o\rangle - \frac{V_{21}}{\Delta} |2^o\rangle$$

$$V_{21} = V_{12}^*, \text{ assume } V_{12} = \text{real} \Rightarrow V_{12} = V_{21} = V$$

$$|2\rangle = |2^o\rangle + \frac{V}{\Delta} |1^o\rangle + |1\rangle = |1^o\rangle - \frac{V}{\Delta} |2^o\rangle$$

$|1\rangle + |2\rangle$  are not normalized

$$\langle 1|1\rangle = \left\{ \langle 1^0| - \frac{v}{\Delta} \langle 2^0| \right\} \left\{ |1^0\rangle - \frac{v}{\Delta} |2^0\rangle \right\}$$

$$= 1 + \frac{v^2}{\Delta^2}$$

$$\langle 2|2\rangle = 1 + v^2/\Delta^2$$

$\Rightarrow$  normalized wavefunctions are

$$|1\rangle = \frac{1}{\sqrt{1+v^2/\Delta^2}} \left\{ |1^0\rangle - \frac{v}{\Delta} |2^0\rangle \right\}$$

$$|2\rangle = \frac{1}{\sqrt{1+v^2/\Delta^2}} \left\{ |2^0\rangle + \frac{v}{\Delta} |1^0\rangle \right\}$$

Note:  $\langle 1|2\rangle = \frac{1}{\sqrt{1+v^2/\Delta^2}} \left\{ \langle 1^0| - \frac{v}{\Delta} \langle 2^0| \right\}$

$$\times \left\{ |2^0\rangle + \frac{v}{\Delta} |1^0\rangle \right\}$$

$$= \frac{1}{\sqrt{1+v^2/\Delta^2}} \left( \frac{v}{\Delta} - \frac{v}{\Delta} \right) = 0$$