

Fourier Optics

start by looking at Fourier Transforms

$$f(t) = \int_{-\infty}^{\infty} F(v) e^{-i2\pi vt} dv$$

$$F(v) = \int_{-\infty}^{\infty} f(t) e^{-i2\pi vt} dt$$

} FT pair

Some properties:

(1) Linear $F_1 + F_2 = \int_{-\infty}^{\infty} (f_1 + f_2) e^{-i2\pi vt} dt$

(2) Translation:

Consider $F(v) = \int f(t) e^{-i2\pi vt} dt$

F.T. of $f(t - \tilde{v})$ is $\int_{-\infty}^{\infty} f(t - \tilde{v}) e^{-i2\pi vt} dt$

write $t' = t - \tilde{v}$ $\Rightarrow t = t' + \tilde{v}$

$$\int f(t - \tilde{v}) e^{-i2\pi vt} dt = \int f(t') e^{-i2\pi v(t' + \tilde{v})} dt'$$

$$= e^{-i2\pi v\tilde{v}} F(v)$$

(2)

\Rightarrow shifting in time by τ introduces
a phase factor $e^{-i2\pi v\tau}$

likewise $\int_{-\infty}^{\infty} F(v-v_0) e^{i2\pi v t} dv$

$$= e^{i2\pi v_0 t} \int F(v') e^{i2\pi v' t} dv'$$

$$= e^{i2\pi v_0 t} f(t)$$

(3) Convolution

~ the convolution of 2 functions $f_1 + f_2$

is $f(t) = \int_{-\infty}^{\infty} f_1(\tau) f_2(t-\tau) d\tau$

in Fourier Space: $F(v) = F_1(v) \times F_2(v)$

$\underbrace{F_1}_{\text{F.T. of } f_1(t)}$ $\overbrace{F_2}^{\text{F.T. of } f_1(t) \text{ and } f_2(t)}$

(3)

Examples:

$f_1(t) = e^{-\delta t}$ \hookrightarrow exponential decay

$f_2(t) = e^{-t^2/\sigma^2}$ \hookrightarrow Gaussian (laser pulse?)

$$F_2(v) = \int_{-\infty}^{\infty} f_2(t) e^{-i2\pi v t} dt$$

$$= \sigma \sqrt{\pi} e^{-\pi^2 \sigma^2 v^2} \quad \hookrightarrow \text{F.T. of a Gaussian is a Gaussian}$$

$$F_1(v) = \int_{-\infty}^{\infty} f_1(t) e^{-i2\pi v t} dt$$

$$= \int_{-\infty}^0 f_1(-t) e^{-i2\pi v t} dt + \int_0^{\infty} f_1(t) e^{-i2\pi v t} dt$$

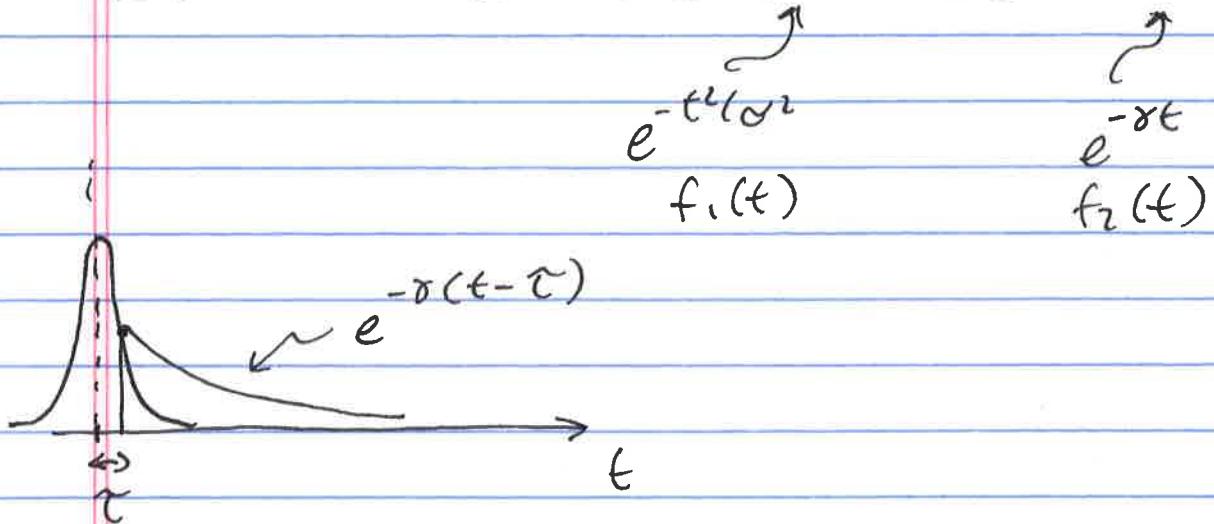
$$= \frac{1}{-2i\pi v^2 + \delta} + \frac{1}{2i\pi v + \delta}$$

$$= \frac{2\delta}{4\pi^2 v^2 + \delta^2} \quad \hookrightarrow \text{F.T. of an exponential decay is a Lorentzian}$$

(4)

Convolution of f_1 and f_2 ?

~ signal in T.R. experiments is a convolution of the IRF and decay

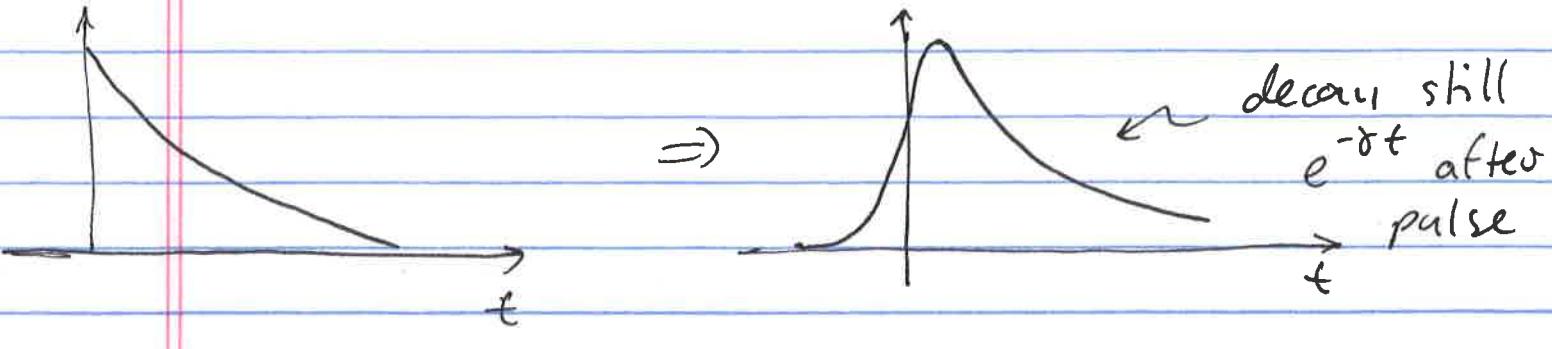


$$\text{Total signal } f(t) = \int_{-\infty}^t f_1(\tau) f_2(t-\tau) d\tau$$

$$= \frac{\alpha \sqrt{\pi}}{2} e^{-\gamma t} \times \left\{ 1 + \operatorname{Erf}(t/\alpha) \right\}$$

"Error Function"

~ not analytic



(5)

~ we can also define spatial Fourier Transforms, i.e.

$$\tilde{F}(v_x) = \int_{-\infty}^{\infty} f(x) e^{-i2\pi(v_x)x} dx$$

$$f(x) = \int_{-\infty}^{\infty} \tilde{F}(v_x) e^{i2\pi(v_x)x} dv_x$$

• spatial frequency $v_x = k_x / 2\pi$

$$2D: F(v_x, v_y) = \iint f(x, y) e^{-i2\pi(v_x)x + i2\pi(v_y)y} dx dy$$

separates into 2 integrals if $f(x, y)$ is separable, otherwise you have "nested" integrals

Convolution is 2D:

$$f(x, y) = \iint f_1(x', y') f_2(x-x', y-y') dx' dy'$$

$$F(v_x, v_y) = F_1(v_x, v_y) F_2(v_x, v_y)$$

(6)

Back to optics...

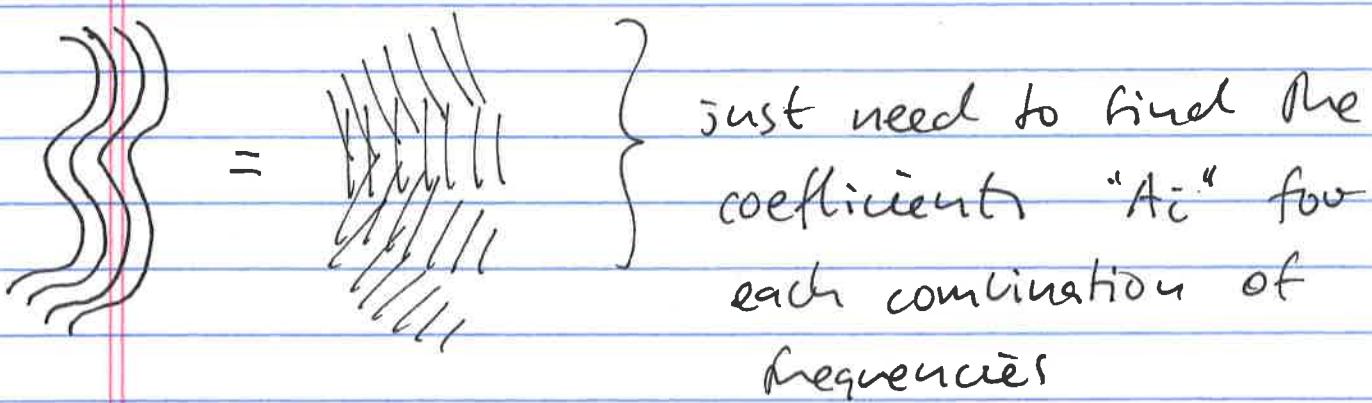
$$K = \frac{2\pi}{\lambda} = \sqrt{k_x^2 + k_y^2 + k_z^2}$$

plane or "harmonic" waves

$$u(x, y, z) = A e^{-i(K_x x + k_y y + k_z z)}$$

$$= A e^{-i2\pi(v_x x + v_y y + v_z z)}$$

~ any wave can be written as a sum of harmonic waves w/ different spatial frequencies



Goal: calculate how a wave changes as it travels through a system or propagates in free space



$$U(x, y, 0) = f(x, y)$$

$$U(x, y, d) = g(x, y)$$

~ we want to find $g(x, y)$ given $f(x, y)$

Note:
$$f(x, y) = \sum_i A_i e^{-i2\pi(v_x^i x + v_y^i y)}$$

$$= \iint F(v_x, v_y) e^{-i2\pi(v_x x + v_y y)} dv_x dv_y$$

$f(x, y) \leftrightarrow F(v_x, v_y) \sim \text{F.T. pair}$

↑
 tells us the amplitude of
 the different spatial frequencies
 that make up $U(x, y, 0)$

~ calculate how $f(x, y)$ is transformed
 into $g(x, y)$ using either impulse-response
 function or transfer function

(8)

~ impulse-response function $h(x, y)$

$$\Rightarrow g(x, y) = \iint_{-\infty}^{\infty} h(x-x', y-y') f(x', y') dx' dy'$$

\uparrow
 $g(x, y)$ is a convolution of $h(x, y) + f(x, y)$

$$\therefore G(v_x, v_y) = H(v_x, v_y) F(v_x, v_y)$$

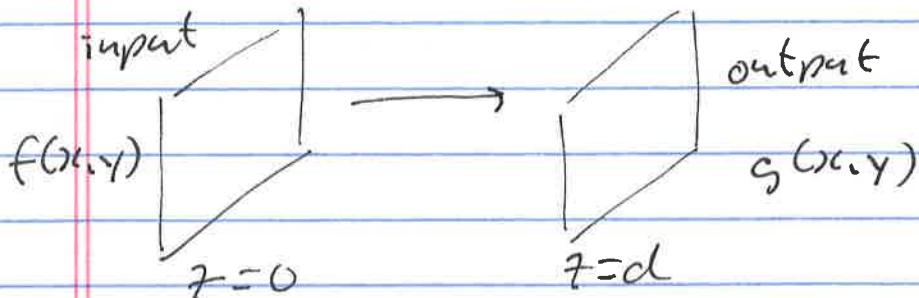
\uparrow
 transfer function

$$= \int h(x, y) e^{-j2\pi(v_x x + v_y y)} dx dy$$

~ for harmonic waves

$$g(x, y) = H(v_x, v_y) f(x, y)$$

~ look at the propagation of
 a plane wave in free space



(a)

$$@ t=0 \quad f(x,y) = u(x,y,0) = A e^{-i 2\pi (v_x x + v_y y)}$$

$$= A e^{-i (k_x x + k_y y)}$$

$$@ t=d \quad g(x,y) = u(x,y,d) = A e^{-i (k_x x + k_y y + k_z d)}$$

where $k_z = (k^2 - k_x^2 - k_y^2)^{1/2}$

$$k = 2\pi/\lambda, \quad k_x = 2\pi v_x, \quad k_y = 2\pi v_y$$

$$\Rightarrow k_z = 2\pi (\frac{1}{\lambda^2} - v_x^2 - v_y^2)^{1/2}$$

Transfer Function: $H(v_x, v_y) = g(x,y)/f(x,y)$

$$= e^{-i k_z d}$$

$$\Rightarrow \left\{ H(v_x, v_y) = e^{-i 2\pi (\frac{1}{\lambda^2} - v_x^2 - v_y^2)^{1/2} d} \right\}$$

2 regimes:

$$(i) \quad v_x^2 + v_y^2 \leq \frac{1}{\lambda^2}$$

$\Rightarrow H(v_x, v_y) \sim$ complex, oscillating function
of d

(10)

$$\Rightarrow |\mathcal{H}(v_x, v_y)|^2 = 1$$

\Rightarrow amplitude doesn't decay

(ii) $v_x^2 + v_y^2 \gg k_\lambda^2$ \sim high spatial frequencies

$$\mathcal{H}(v_x, v_y) = e^{-iz\tilde{r}(k_\lambda^2 - v_x^2 - v_y^2)^{1/2}i}$$

imaginary

$\Rightarrow \mathcal{H}(v_x, v_y)$ decays w/ distance

"evanescent" waves

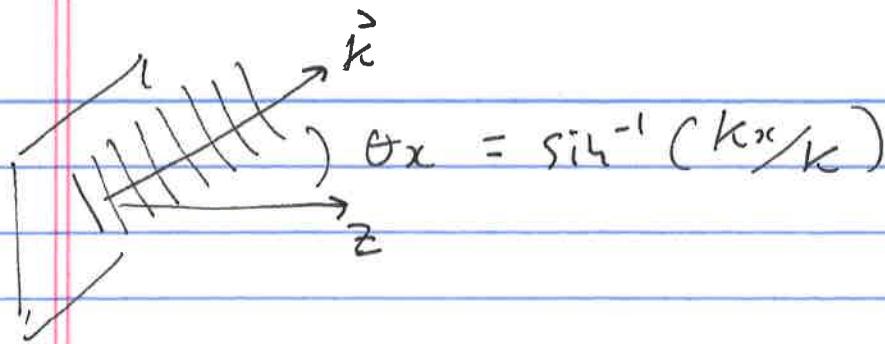
\sim high spatial frequencies \Leftrightarrow small distances

\Rightarrow if we could collect high v components
we could improve spatial resolution
(e.g. NSOM)

Fresnel Approximation:

$$v_x^2 + v_y^2 \ll k_\lambda^2$$

\sim waves that make small angles
w.r.t z -axis



$$\theta_x = \sin^{-1}(k_x/k)$$

~ small angles $\theta_x \approx k_x/k = \frac{k_x d}{\lambda \pi} = v_x \lambda \ll 1$

"Paraxial Rays"

$$\Rightarrow 2\pi \left(\frac{1}{\lambda^2} - v_x^2 - v_y^2 \right)^{1/2} d = \frac{2\pi d}{\lambda} (1 - \theta^2)^{1/2}$$

where $\theta^2 = \lambda^2(v_x^2 + v_y^2) \ll 1$

$$\Rightarrow \frac{2\pi d}{\lambda} (1 - \theta^2)^{1/2} \approx \frac{2\pi d}{\lambda} (1 - \theta^2/2)$$

and $H(v_x, v_y) = e^{-i \frac{2\pi d}{\lambda} (1 - \theta^2/2)}$

$$= H_0 e^{+i \frac{\pi d}{\lambda} \theta^2}$$

$$= H_0 e^{i \pi \lambda d (v_x^2 + v_y^2)}$$

$$= H_0 e^{\uparrow}$$

$$e^{-i \frac{2\pi d}{\lambda} \theta^2} = e^{-ikd}$$

Transfer function of free space, in the Fresnel approximation, is

$$H(v_x, v_y) = H_0 e^{i \pi d (v_x^2 + v_y^2)}$$

impulse-response function

$$\begin{aligned} h(x, y) &= \iint H(v_x, v_y) e^{-i \pi (v_x x + v_y y)} dv_x dv_y \\ &= h_0 e^{-ik(x^2 + y^2)/2d}, \\ h_0 &= \frac{i}{\lambda d} e^{-ikd} \end{aligned}$$

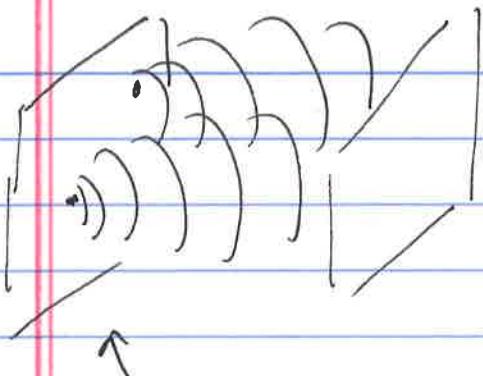
"Paraboloid" wave

rewrite an expression for $g(x, y)$ using $h(x, y)$

$$g(x, y) = \iint f(x', y') h(x - x', y - y') dx' dy'$$



Huygen's Principle: each pt. on the input plane generates a spherical wave, we add them all up at the output plane to get the total field



Fresnel approx \approx treat spherical waves (e^{-ikr}/r) as paraboloid waves
 $e^{-i\pi(x^2+y^2)/\lambda d}/d$

Example: propagation of a Gaussian

$$f(x, y) = A e^{-((x^2+y^2)/w_0^2)}$$

$$g(x, y) = h_0 \iint f(x', y') e^{-i\frac{k((x-x')^2+(y-y')^2)}{2z}} dx' dy'$$

$$h_0 = \frac{i}{\lambda z} e^{-ikz}$$

$$\iint f(x') e^{-i\frac{k((x-x')^2)}{2z}} dx' = \frac{\sqrt{2\pi} e^{-kx^2/(kw_0^2+2iz)}}{\sqrt{\frac{2}{w_0^2} + \frac{ik}{z}}}$$

Integral over y gives same term

14

$$\frac{i}{z} \frac{2\pi}{(\frac{2}{\omega_0^2} + \frac{cK}{z})} e^{-(x^2+y^2)} \frac{K}{kw_0^2 - 2iz} e^{-ikz}$$

$$z_0 = Kw_0^2/2$$

$$\frac{ik}{(\frac{2z_0}{\omega_0^2} + ik)} = \frac{-i(Kw_0^2/2)}{(z + ikw_0^2/2)} = \frac{cz_0}{(z + cz_0)}$$

cf 3.1-5

argument of exponential

$$\frac{-\rho^2 k}{kw_0^2 - 2iz} = \frac{-\rho^2 k}{k^2 w_0^4 + 4z^2} \times (kw_0^2 + z - z)$$

Real part: $\frac{-\rho^2 k^2 w_0^2}{k^2 w_0^4 + 4z^2} = \frac{-\rho^2}{w_0^2 (1 + \frac{4z^2}{k^2 w_0^4})}$

$$= \frac{-\rho^2}{w_0^2 (1 + (z/z_0)^2)}$$

Imaginary part: $\frac{-\rho^2 2izk}{k^2 w_0^4 + 4z^2} = \frac{-ik\rho^2}{2z \left(\frac{k^2 w_0^4}{4z^2} + 1 \right)}$

$$= -ik\rho^2 \cancel{\frac{1}{2z \left(1 + (\frac{z_0}{z})^2 \right)}}$$

~ expression derived from consideration
 $w(h(x,y))$ is equivalent to the
 Gaussian beam formula

i.e. $\frac{i z_0}{z + i z_0} e^{-\rho^2/\omega(z)^2} e^{-ikz - ik\rho^2/2R(z)}$

where $\omega(z) = \omega_0 (1 + (z/z_0)^2)^{1/2}$

$R(z) = z (1 + (\tau_0(z))^2)$